

Analyzing the effects of post couplers in DTL tuning by the equivalent circuit model*

JIA Xiao-Yu(贾晓宇)^{1,2} ZHENG Shu-Xin(郑曙昕)^{1,2;1)}

¹ Key Laboratory of Particle & Radiation Imaging (Tsinghua University), Ministry of Education, Beijing 100084, China

² Department of Engineering Physics, Tsinghua University, Beijing 100084, China

Abstract: Stabilization of the accelerating field in Drift Tube Linac(DTL) is obtained by inserting Post Couplers(PCs). On the basis of the equivalent circuit model for the DTL with and without asymmetrical PCs, stabilization is deduced quantitatively: we let $\delta\omega/\omega_0$ be the relative frequency error, then we discover that the sensitivity of field to perturbation is proportional to $\sqrt{\delta\omega/\omega_0}$ without PCs and to $\delta\omega/\omega_0$ with PCs. Then we adapt the circuit model of symmetrical PCs for the case of asymmetrical PCs. The circuit model shows how the slope of field distribution is changed by rotating the asymmetrical PCs and illustrates that the asymmetrical PCs have the same effect as the symmetrical ones in stabilization.

Key words: equivalent circuit model, DTL tuning, stabilization

PACS: 29.20.Ej **DOI:** 10.1088/1674-1137/37/12/127005

1 Introduction

DTL is a kind of standing-wave structure and operates at TM₀₁₀ mode. In this mode, the slope of dispersion curve is zero, which means the field distribution of this structure is sensitive to perturbations, such as beam loading, machining and installation errors [1–3]. We can reduce the effect of the perturbations by inserting the PCs, which are used to introduce resonant coupling mode and increase the slope of dispersion curve at the operating point [4]. The tuning for DTL consists of the following tasks: resonant frequency of the structure, stabilization and distribution of accelerating electric field. Generally speaking, we can adjust slag tuners to get resonant frequency, insert the PCs into the right length to improve the stability of the field and rotate the PCs to adjust the field distribution precisely [5].

In order to reduce the interaction, the adjacent PCs are on different sides of the drift tube and all PCs are perpendicular to the stem. Both symmetrical and asymmetrical PCs can make the field stable. During the stabilization state, field distribution is flat when using symmetrical PCs, but we can get a tilted field distribution with asymmetrical ones and the rate of slope can be changed by rotating the PCs [6]. Fig. 1 shows the structure of DTL and the asymmetrical PCs.

References [7–9] propose an equivalent circuit model for the DTL with the symmetrical PCs; further, this arti-

cle improves the model so that field distribution and stabilization with asymmetrical PCs will be analyzed and the rate of the tilting field can be estimated.

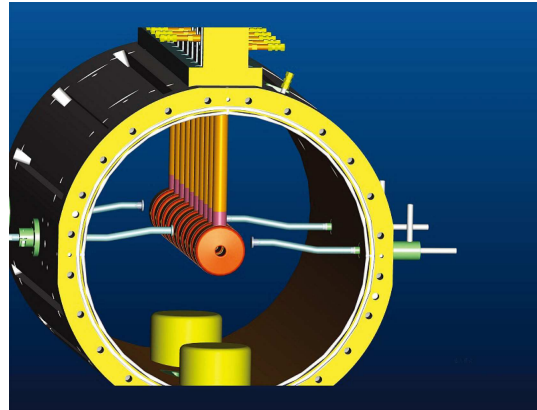


Fig. 1. DTL and the post couplers (PCs).

2 The equivalent circuit for general periodic structure

General periodic structure can be represented by the circuit in Fig. 2, which is composed of equivalent parallel admittance Y_n and series impedance Z_n . By selecting the loop with Y_{n-1} , Z_n and Y_n , we can get the recurrence formula of current I_n .

$$-(I_{n-1} - I_n)/Y_{n-1} + I_n Z_n + (I_n - I_{n+1})/Y_n = 0. \quad (1)$$

Received 28 March 2013

* Supported by National Natural Science Foundation of China (91126003)

1) E-mail: zhengsx8@gmail.com

©2013 Chinese Physical Society and the Institute of High Energy Physics of the Chinese Academy of Sciences and the Institute of Modern Physics of the Chinese Academy of Sciences and IOP Publishing Ltd

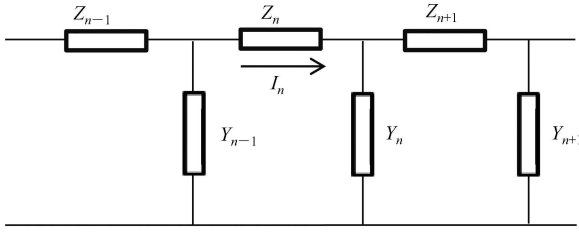


Fig. 2. The equivalent circuit for DTL without PCs.

The character of the periodic structure means $Y_{n-1} = Y_n = Y$, $Z_n = Z$, then the recurrence formula can be simplified.

$$\begin{cases} I_{n+1} - (2+YZ)I_n + I_{n-1} = 0 & Y \neq 0. \\ I_{n+1} = I_n = I_{n-1} & Y = 0. \end{cases} \quad (2)$$

Y can not be equal to 0 in Eq. (1), and the situation of $Y = 0$ is considered in Eq. (2). There are different solutions for other situations of YZ .

1) $Y=0$: from Eq. (2) we know that $I_{n+1} = I_n = I_{n-1}$. It is open circuit state and current through every cell is equal to each other. We call this situation the flat-mode.

2) $Z=0$ and $Y \neq 0$: the solution is $I_n = I_{n-1} + \Delta I$. Current to the next cell will be added by the same ΔI , which is the current through admittance Y . Now the current distribution is an oblique line and we call this situation the tilt-mode.

3) $YZ \neq 0$: the solution is

$$I_n = I_{n-1} e^{\pm i\phi}, \quad \cos\phi = 1 + YZ/2. \quad (3)$$

YZ can be complex and ϕ also can be complex. This is the most general solution, but in an ideal state in DTL, power loss is not considered, and Y, Z are near to 0 (that will be explained in Chapter 3). Then we will concentrate on several situations closely related to tuning.

4) $|YZ+2| < 2$ and YZ is real: the solution is

$$I_n = I_{n-1} e^{\pm i\phi}, \quad \cos\phi = 1 + YZ/2. \quad (4)$$

This is a special situation of Case 3). Here ϕ is real, and it is the phase difference between adjacent cells. This solution represents the transmission electromagnetic wave and ϕ is the phase shift between adjacent cells. We call this situation the transmission-mode.

5) $|YZ+2| > 2$ and YZ is real: the solution is

$$I_n = I_{n-1} r, \quad r = \frac{2 + YZ \pm \sqrt{(2 + YZ)^2 - 4}}{2}. \quad (5)$$

It is also a special situation exception of Case 3). The amplitude of current attenuates according to the exponential rate, r is the current ratio of adjacent cell. In DTL, this solution means the electromagnetic wave cuts off, and the amplitude of the field attenuates according to the exponential rate. But if r is close to 1, the field distribution would be similar to an oblique line within suitable range. We call this situation the cutoff-mode.

Table 1 summarizes the important solutions.

As is known, Y and Z are as functions of the frequency ω , so different frequencies correspond to different modes. If the RF source frequency differs from the structure resonant frequency by $\delta\omega$, then the field distribution would deviate from the design value. For the cutoff-mode, stabilization can be measured by the r , which is the current ratio of adjacent cell, and for the transmission-mode, it can be measured by the phase change of ϕ . Then we will prove this change is proportional to $\delta\omega/\omega_0$ with PCs in the right position, and proportional to $\sqrt{\delta\omega/\omega_0}$ without PCs. For example, the operating frequency is 325 MHz, and the difference between the RF source frequency and the structure resonant frequency is less than 30 kHz. In this case, the field distribution would change by about 0.01% with PCs and about 1% without PCs. We can see that PCs can improve the stabilization significantly.

3 The equivalent circuit for DTL

DTL is a kind of quasi-periodic structure, but in this paper we will approximately regard it as a periodic one. The power loss is ignored, so there is no resistance in the circuit.

3.1 The equivalent circuit for DTL without PCs

In the circuit model in Fig. 3, C_0 is the capacitance between two drift tubes; L_0 is the inductance of the drift tube; C_s is the capacitance between drift tube and tank; L_s is the inductance of the stem. There are two resonant units in the circuit: one consists of L_0 and C_0 and the other is made up by L_s and C_s . Define $\omega_0 = 1/\sqrt{L_0 C_0}$

Table 1. Solutions for Eq. (2).

condition	name for solution	current relation
$Z=0$ and $Y \neq 0$	tiltmode	$I_n = I_{n-1} + \Delta I$
$Y=0$	platmode	$I_n = I_{n-1}$
$ YZ+2 < 2$ and YZ is real	transmissionmode	$I_n = I_{n-1} e^{\pm i\phi}, \cos\phi = 1 + \frac{YZ}{2}$
$ YZ+2 > 2$ and YZ is real	cutoffmode	$I_n = I_{n-1} r, r = \frac{2 + YZ \pm \sqrt{(2 + YZ)^2 - 4}}{2}$

and $\omega_s = 1/\sqrt{L_s C_s}$. As we know, ω_s is much lower than ω_0 , the operating frequency. So we can regard L_s and C_s as a capacitance C parallel in the circuit [7].

$$C = C_s \frac{\omega^2 - \omega_s^2}{\omega^2}. \quad (6)$$

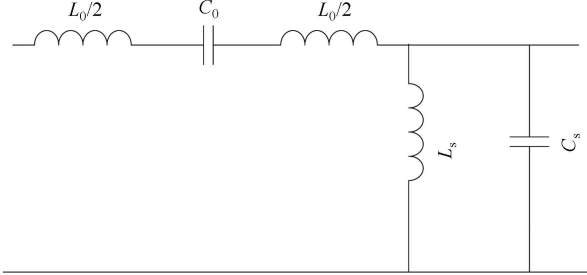


Fig. 3. The equivalent circuit for DTL without PCs.

Voltage of the capacitance C_0 can be considered as the product of cell length and average accelerating field: $u_n = L_{\text{cell}} E_{n,\text{ave}}$, so the current i_n through C_0 is related to the field: $i_n = j\omega C_0 u_n = j\omega C_0 L_{\text{cell}} E_{n,\text{ave}}$.

The operating frequency is equal to the resonant frequency of DTL, which means $\omega = \omega_0$, so $Z=0$, $Y = j\omega_0 C$. According to Table 1, it is the tilt-mode and the slope of field is decided by boundary condition, which is the end cavities frequency of DTL. In practice, we can tilt the field by changing the frequency of the end cavities. When the operating frequency is shifted by $\delta\omega$, by linear approximation we can get

$$YZ = \frac{C}{C_0} \frac{\omega_0^2 - (\omega_0 + \delta\omega)^2}{\omega^2} \approx -2 \frac{C}{C_0} \frac{\delta\omega}{\omega_0}. \quad (7)$$

For $\delta\omega/\omega \ll 1$, so $|YZ| \ll 1$, if $\delta\omega > 0$, then $YZ < 0$, according to Table 1, the field would stay in transmission-mode and

$$\Delta\phi \approx \sqrt{2C\delta\omega/(C_0\omega_0)};$$

$\delta\omega > 0$, the field is in cutoff-mode and

$$\Delta r \approx \pm \sqrt{2C\delta\omega/(C_0\omega_0)}.$$

Because C/C_0 is constant, the sensitivity of the field to perturbation is in proportion to $\sqrt{\delta\omega/\omega_0}$, and the stabilization is poor.

3.2 The equivalent circuit for DTL with symmetrical PCs

C_p in Fig. 4 is the capacitance between PC and drift tube and L_p is the inductance of PC, then we define $\omega_p = 1/\sqrt{L_p C_p}$. By inserting PCs, we introduce a new resonant mode, and the stabilization can be acquired if this mode is completely coupled with the operation mode. Coupling condition can be proposed by solving

the dispersion equation [1]

$$\cos\phi = 1 + \frac{Y(\omega)Z(\omega)}{2}. \quad (8)$$

It exactly means the parallel admittance $Y=0$. Corresponding to Table 1, it is the flat-mode and the distribution of field is unrelated to boundary conditions.

$$\omega_p^2 = \frac{C}{C+C_p} \omega_0^2. \quad (9)$$

When the difference between the RF source frequency and the structure resonant frequency is $\delta\omega$, by linear approximation we can get

$$YZ \approx -4 \frac{C(C+C_p)}{C_0 C_p} \frac{\delta\omega^2}{\omega_0^2}. \quad (10)$$

$|YZ| \ll 1$ and $YZ < 0$, so the structure operates in the transmission-mode and

$$\Delta\phi \approx \frac{\delta\omega}{\omega} \sqrt{\frac{C(C+C_p)}{C_0 C_p}}. \quad (11)$$

The sensitivity of field to perturbation is in proportion to $\delta\omega/\omega$, and the stabilization is improved compared with that without PCs.

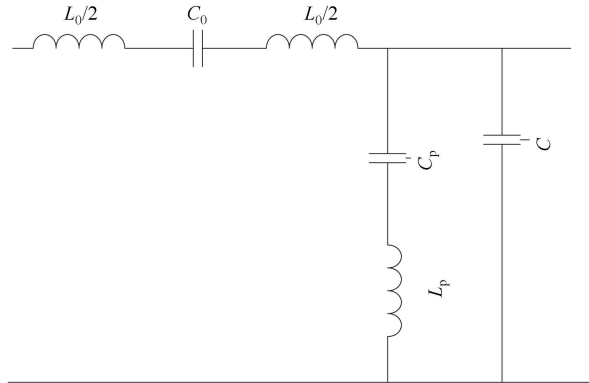


Fig. 4. The equivalent circuit for DTL without PCs.

3.3 The equivalent circuit for DTL with asymmetrical PCs

While rotated, the tip position of asymmetrical PCs changes, so the positions of C_p and L_p in circuit change at the same time, as Fig. 5(b). Δ -Y transform is taken for the box in Fig. 6(b), so we can get a simple circuit with the same form of Fig. 2. Because of $\delta L/L_0 \ll 1$, when PCs are tuned to satisfy Eq. (9), we can get

$$YZ \approx \left(\frac{C\delta L}{C_0 L_0} \right)^2. \quad (12)$$

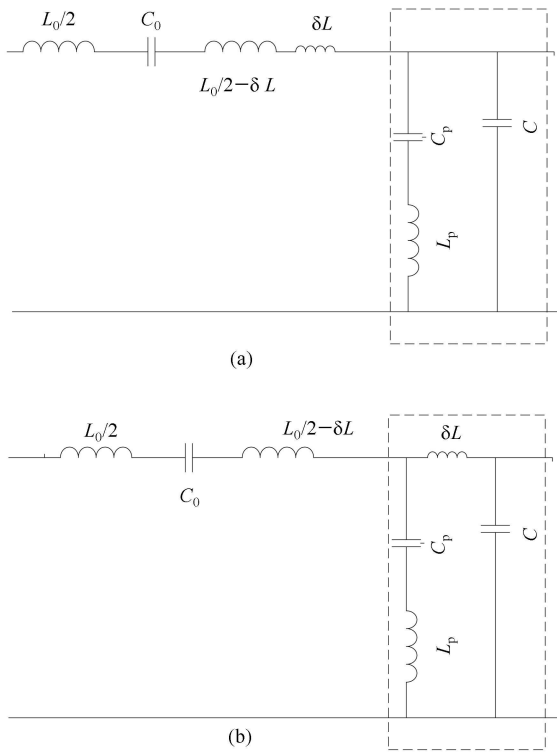


Fig. 5. The circuit for DTL with symmetrical PCs (a) and with asymmetrical PCs (b).

The structure works in the cutoff-mode, and $r \approx \pm C\delta L/(C_0L_0)$. Because r (defined in Chapter 2) is quite close to 1, which means the distribution of the current would look like a tilt line whose slope would be affected by the end cavities and the value of δL . The slope of the amplitude of the electric field can be adjusted by varying the value of δL , that means rotating PCs. Then we analyze the stabilization with the above method. If

$$\frac{\delta\omega}{\omega} \ll 1, \frac{\delta L}{L_0} \ll 1, \frac{\delta\omega}{\omega} / \frac{\delta L}{L_0} \ll 1, \quad (13)$$

is satisfied, the result will be obvious. $\Delta r \approx \pm C\delta\omega/(C_0\omega_0)$, it illustrates that asymmetrical PCs have the same effect as the symmetrical ones in stabilization.

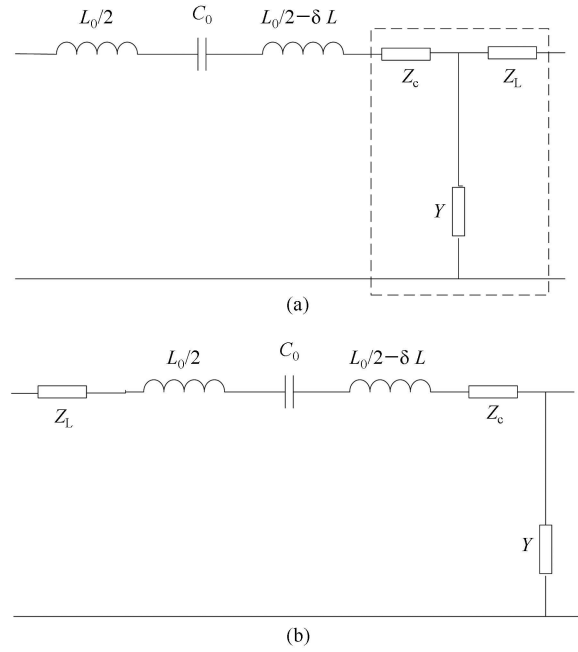


Fig. 6. The result for Δ -Y transform.

4 Conclusion

On the basis of the equivalent circuit model, the character of the stabilization and distribution of the electric field in DTL was analyzed in this article. Firstly, the stabilization was described by the relationship between field perturbation and frequency difference: the sensitivity of field to perturbation is in proportion to $\sqrt{\delta\omega/\omega_0}$ without PCs and to $\delta\omega/\omega_0$ with PCs. Secondly, the difference between the symmetrical and asymmetrical PCs is that the electric field distribution is flat and unrelated to the adjustment of the end cavities with the symmetrical PCs, but that attenuates according to exponential rate with asymmetrical PCs. If r is close to 1, the field distribution can be considered as an oblique line and the slope of it can be changed by rotating the asymmetrical PCs.

References

- 1 Wangler T. RF Linear Accelerators. Second Edition. New York: John Wiley & Sons Publishing Company, Inc., 2008. 57–72
- 2 Naito F, Tanaka H, Ikegami M et al. Tuning of the RF Field of the DTL for the J-Parc. In: Proc. of PAC 2003. Conf. on Particle Accelerator. IEEE, 2003, **5**: 2835
- 3 Billen J, Garcia J, Potter J et al. Nuclear Science, 1985, **32**(5): 3184
- 4 Knapp E. US Patent 3, 501, 734. 1970
- 5 Machida S, Kato T, Fukumoto S. Nuclear Science, 1985, **32**(5): 3259
- 6 Roy S, Pande R, Rao S et al. Electromagnetic Design of DTL Cavity for Lehipa. In: Proc. DAE-BRNS Symposium on Ion Beam Technology and Applications, 2007. 333
- 7 Grespan F, de Michele G, Ramberger S et al. Circuitual Model for Post Coupler Stabilization in a Drift Tube Linac. Tech. Rep., 2010, sLHC Project Note 0014
- 8 Vretenar M. arXiv: 1201.2593
- 9 Ke M, Cheo B, Shmoys J. Particle Accelerators, 1994, **48**(2): 109