

EC/ β^+ decay of six medium-heavy nuclei

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Abstract Previous experimental results of (EC+ β^+) decay for the medium-heavy nuclei reported by our group since 1996, including ^{153}Er , ^{157}Yb , ^{209}Fr , ^{128}Ce , ^{130}Ce , and ^{128}Pr have been briefly summarized. The observed low-lying states in their daughter nuclei have been reviewed in a systematic way and compared with different model calculations. Finally, some questions have been put forward for further study and discussion.

Key words (EC+ β^+) decay, low-lying states, nuclear shape, multiplet, Gamow-Teller transition

PACS 23.40.-s, 27.90.+b

1 Introduction

Study of EC/ β^+ decay is an important approach to investigate the nuclear structure in low-lying region. The Gamow-Teller transition is the major decay mode of neutron-deficient medium-heavy nuclei. In medium-heavy mass region, especially in the rare-earth region the nuclear shape changes rather complicatedly. Therefore, proposed EC/ β^+ decay scheme could provide important information not only for nuclear shape and related structure in the region, but also for Gamow-Teller transition itself.

Our group has observed EC/ β^+ decays for ^{153}Er , ^{157}Yb , ^{209}Fr , ^{128}Ce , ^{130}Ce , and ^{128}Pr , and proposed decay schemes for all of them^[1–5] since 1996. The aimed nuclei were produced via the fusion evaporation reactions induced by heavy ions ^{16}O and ^{36}Ar , which were provided from a Sector Focused Cyclotron at the Institute of Modern Physics, Lanzhou, China. We used a He-jet fast tape transport system to move the reaction products to a shielded counting room, where the X- γ - γ - t coincidence measurements were carried out. In addition, the X- γ coincidence measurements were also used for Z identification. Sometimes, the excitation-function measurements were used for mass identification. In this paper we simply summarize the proposed decay schemes, make a systematic review to the observed low-lying states in the daughter nuclei, and compare the ob-

tained physical results with different model calculations. Finally, some questions are put forward for further study and discussion.

2 Results and discussion

2.1 Even(Z)-odd(N) nuclei

From the EC/ β^+ decay scheme of ^{153}Er , we found single particle states $s_{1/2}$, $d_{3/2}$ and maybe $d_{5/2}$ as well as a three quasi-particle state in the low-lying region of the daughter nucleus ^{153}Ho , while in the EC/ β^+ decay scheme of ^{157}Yb we found a rotational band in the low-lying region of the daughter nucleus ^{157}Tm ^[1]. The Fig. 4 of the Ref. [1] is the systematic behavior of some characteristic low-lying states in the odd- A Tm($Z=69$) and odd- A Ho($Z=67$) isotopic chains with $N = 82–90$, and shows that the single particle states dominant the low-lying region in the isotopes with $N = 82–86$, while the rotational bands dominant the low-lying region in the isotopes with $N = 88$ and 90 .

This fact indicates that the transition point from near-spherical ground state to deformed ground-state appears between $N = 86$ and 88 in the two odd- A isotopic chains. However, as J. H. Hamilton^[6] pointed out, a sudden onset of deformation happens between $N = 88$ and 90 in the even-even isotopic chains around the weak $Z = 64$ spherical shell region,

Received 3 September 2008

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such as Sm($Z=62$), Gd($Z=64$) and Dy($Z=66$). The two-neutron shift for the transition point from Odd- A isotopic chain to Even-Even isotopic chain probably due to the paring effect which tends to maintain the spherical shape of nuclear ground state. The predicted deformation as a function of neutron number for the element Dy, Ho and Er is shown in Fig. 1. It can be seen in Fig. 1 that a sudden change of the nuclear deformation happens between $N = 84$ and 86 . The theoretical predictions were given by Möller et al based on their macroscopic-microscopic model^[7], and not consistent with both experimental results. So far we do not know what the reason behind the inconsistency is.

2.2 Odd(Z)-even(N) nucleus

The EC/ β^+ decay scheme of ^{209}Fr was proposed by our group^[2] in 1996. The branching ratio of the EC/ β^+ decay of ^{209}Fr is as weak as 3%. Based on the decay scheme we found a multiplet, i.e. a five-fold state [$(^{210}\text{Rn } 2^+)(\nu f_{5/2})^{-1}$] in the low-lying region of the daughter nucleus of ^{209}Rn . Here, the ($^{210}\text{Rn } 2^+$) stands for the 2^+ vibration state of the core ^{210}Rn , while the $(\nu f_{5/2})^{-1}$ stands for the neu-

tron hole at the orbital $f_{5/2}$. There are 123 neutrons in ^{209}Rn . The last neutron in ^{209}Rn locates at the orbital $5/2[503]$. The analogical structure of low-lying states in the $N=123$ isotone chain is shown in Fig. 2. The similar multiplet, which consists of a 2^+ vibrational state of the core and the neutron hole at the orbital of $f_{5/2}$, can be seen at the low-lying region not only in $^{209}\text{Rn}(Z=86)$, but also in $^{207}\text{Po}(Z=84)$ and $^{205}\text{Pb}(Z=82)$. However, we are not able to reproduce the systematics by a shell model calculation yet.

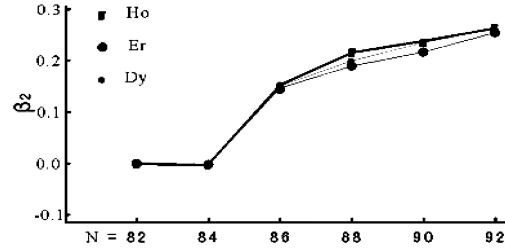


Fig. 1. Predicted deformation β_2 as a function of neutron number for the element Dy, Ho and Er.

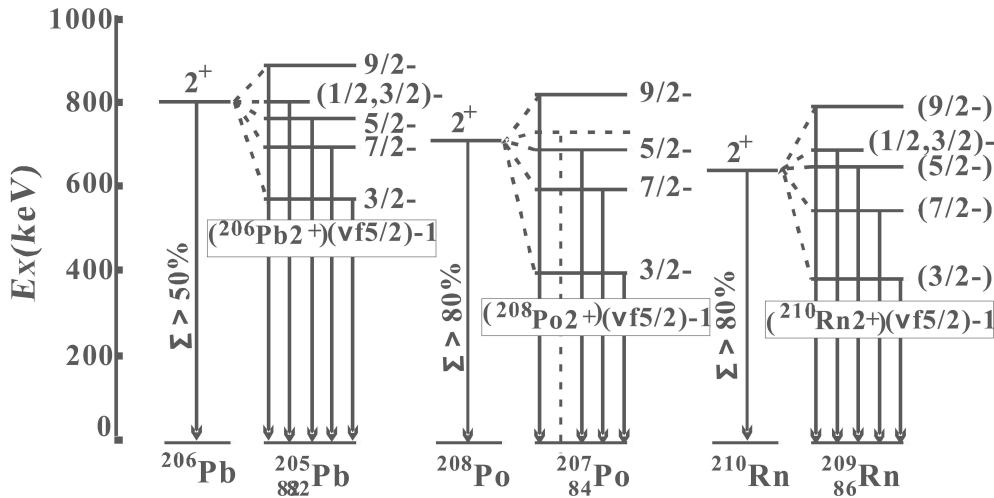


Fig. 2. Analogical Structure of low-lying states in the isotone chain with $N = 123$.

2.3 Even-even nuclei

The EC/ β^+ decay scheme of ^{130}Ce was proposed by our group in 1996^[3], and evaluated and edited in Nuclear Data Sheets by Singh et al in 2001^[8]. The experiment was carried out very carefully. The energy of γ line was searched up to Q_{EC} value. It means that the Q_{EC} window was covered completely. Furthermore, the intensity ratio for the observed most intense γ line over the observed weakest γ line reached

to 5×10^3 . In the decay scheme we found 13 1^+ states in low-lying region of the daughter nucleus ^{130}La . 13 1^+ states were also found in the low-lying region of ^{128}La via the EC/ β^+ decay scheme of ^{128}Ce reported by our group in 1999^[4]. The low-lying 1^+ states in La populated by EC/ β^+ decay of Ce is shown in Fig. 3. The experimental data for $A = 132$ and 134 were reported by Abdurazakov^[9] and Islamov^[10], respectively. Using QPNM (quasi-particle-phonon nuclear model) and GTpp interaction, i. e. a RPA calculation

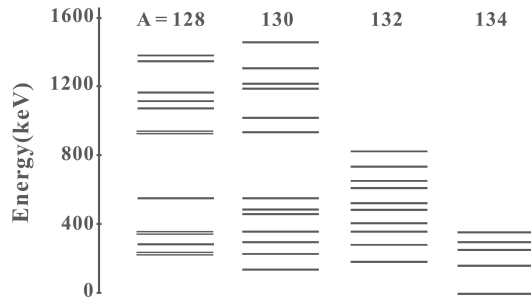


Fig. 3. Low-lying 1^+ states in La isotopes populated by EC/β^+ decay of even-even Ce isotopes with $A = 128$ to 134 .

including separable p-p and p-h interaction, Kuzmin & Soloviev^[11] could reproduce Gamow-Teller strength for the EC/β^+ decay near the double magic nuclei, such as ^{100}Sn and ^{146}Gd . They made a similar calculation for the EC/β^+ decay of ^{130}Ce based on the spin-flip mode $\pi d_{5/2} \rightarrow \nu d_{3/2}$. The calculated $\log ft$ value, 4.2, could be consistent with the experimental integrated $\log ft$ value, i. e. the $\log ft$ value integrated over the 13 1^+ states in ^{130}La . We don't know why the single particle configuration $(\pi d_{5/2}, \nu d_{2/3})1^+$

fragments, and then the Gamow-Teller transition is fed into 13 1^+ states. What interaction leads to the fragmentation?

2.4 Odd-odd nucleus

If the ground-state band in the daughter, even-even nucleus has been already known, the ground-state spin and parity for the mother, odd-odd nucleus can be determined by its EC/β^+ decay. According to $\log ft$ value, the allowed transitions are fed to 2^+ , 4^+ and 3^+ states in ^{128}Ce in the proposed EC/β^+ decay scheme of ^{128}Pr ^[5]. Therefore the spin and parity of ^{128}Pr ground state was assigned to be 3^+ , which is not consistent with the theoretical predictions 5^+ with the configuration $\nu 7/2[523] \times \pi 3/2[541]$, given by Möller et al based on their macroscopic-microscopic model^[12]. However, the projected shell model calculation^[5] could reproduce the experimental assignments 3^+ with the configuration $\nu 1/2[541] \times \pi 5/2[532]$ and a large quadrupole deformation $\beta_2 = 0.408$. Is the large deformation reasonable?

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