

相互作用玻色子-费米子模型的 Spin(6)的极限

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摘 要

本文以李群的表示理论为基础对原子核的相互作用玻色子-费米子模型 (IBFM II) 的 Spin(6) 极限情况作进一步探讨, 用旋量与非旋量外积扬 (Young) 图法给出了群链

$$U_{\pi}^{(B)}(6) \otimes U_{\nu}^{(B)}(6) \otimes U^{(F)}(4) \supset SO_{\pi}^{(B)}(6) \otimes SO_{\nu}^{(B)}(6) \otimes SU^{(F)}(4) \\ \supset SO_{\pi+\nu}^{(B)}(6) \otimes SU^{(F)}(4) \supset Spin(6) \supset Spin(5) \supset Spin(3)$$

的有关约化规则, 按这一群链分类的波函数以及具有这种动力学对称性的能谱。

一、引 言

相互作用玻色子模型 (IBM) 已成为核结构理论研究的重要课题, 相互作用玻色子-费米子模型 (IBFM) 则是 IBM 的进一步发展。

F. Iachello 和 S. Kuyucak 对 IBFM(I) 的 Spin(6) 极限情况作了初步讨论, 他们所考虑的原子核由 n 个 s, d 玻色子和一个 $j = 3/2$ 的费米子构成。

我们所考虑的原子核是由 n_{ν} 和 n_{π} 个 (s_{π}, d_{π} 和 s_{ν}, d_{ν}) 玻色子和一个费米子 ($j = 3/2$) 构成。我们用文献[2]旋量与非旋量外积扬 (Young) 图法给出了群链

$$U_{\pi}^{(B)}(6) \otimes U_{\nu}^{(B)}(6) \otimes U^{(F)}(4) \supset SO_{\pi}^{(B)}(6) \otimes SO_{\nu}^{(B)}(6) \otimes SU^{(F)}(4) \\ \supset SO_{\pi+\nu}^{(B)}(6) \otimes SU^{(F)}(4) \supset Spin(6) \supset Spin(5) \supset Spin(3). \quad (1.1)$$

的有关约化规则, 并且给出了按这一链分类的波函数以及具有这种动力学对称性的能谱。

二、各群链的约化

(1) $U^{(B)}(6)$ 到 $SO^{(B)}(6)$ 的约化。这已由文献[3]给出, 其中 $U^{(B)}(6) = U_{\pi}^{(B)}(6) \otimes U_{\nu}^{(B)}(6)$ 和 $SO_{\pi}^{(B)}(6) \otimes SO_{\nu}^{(B)}(6) = SO_{\pi+\nu}^{(B)}(6)$, 并注意到这一约化不是完全可约的。

(2) $SO_{\pi+\nu}^{(B)}(6) \otimes SU^{(F)}(4)$ 到 Spin(6) 的约化。利用旋量和非旋量外积扬 (Young) 图法^[2], 我们导出下列约化规则。

$$(\sigma 0 0) \otimes \left(\frac{1}{2} \frac{1}{2} \frac{1}{2} \right) = (\sigma + 1/2 \ 1/2 \ 1/2) + (\sigma - 1/2 \ 1/2 \ 1/2) \quad (2.1)$$

$$\begin{aligned} (\sigma 0 0) \otimes (\sigma' + 1/2 \ 1/2 \ 1/2) &= (\sigma + \sigma' + 1/2 \ 1/2 \ 1/2) \\ &+ (\sigma + \sigma' - 1/2 \ 3/2 \ 1/2) + (\sigma - \sigma' - 1/2 \ 1/2 \ 1/2) \\ &+ (\sigma - \sigma' - 3/2 \ 5/2 \ 1/2) + (\sigma + \sigma' - 3/2 \ 3/2 \ 1/2) \\ &+ (\sigma + \sigma' - 3/2 \ 1/2 \ 1/2) + \cdots + (\sigma + 1/2 \ \sigma' \\ &+ 1/2 \ 1/2) + (\sigma + 1/2 \ \sigma' - 1/2 \ 1/2) + \cdots \\ &+ (\sigma + 1/2 \ 1/2 \ 1/2) + (\sigma - 1/2 \ \sigma' + 1/2 \ 1/2) \\ &+ (\sigma - 1/2 \ \sigma' - 1/2 \ 1/2) + \cdots + (\sigma - 1/2 \ 1/2 \ 1/2) \\ &+ (\sigma - 3/2 \ \sigma' - 1/2 \ 1/2) + (\sigma - 3/2 \ \sigma' - 3/2 \ 1/2) \\ &+ \cdots + (\sigma - 3/2 \ 1/2 \ 1/2) + \cdots + (\sigma - \sigma' \\ &+ 1/2 \ 3/2 \ 1/2) + (\sigma - \sigma' + 1/2 \ 1/2 \ 1/2) + (\sigma - \sigma' \\ &- 1/2 \ 1/2 \ 1/2) \quad \sigma > \sigma', \end{aligned} \quad (2.2)$$

$$\begin{aligned} (\sigma 0 0) \otimes (\sigma + 1/2 \ 1/2 \ 1/2) &= (2\sigma + 1/2 \ 1/2 \ 1/2) + (2\sigma \\ &- 1/2 \ 3/2 \ 1/2) + (2\sigma - 1/2 \ 1/2 \ 1/2) + (2\sigma - 3/2 \ 5/2 \ 1/2) \\ &+ (2\sigma - 3/2 \ 3/2 \ 1/2) + (2\sigma - 3/2 \ 1/2 \ 1/2) + \cdots \\ &+ (\sigma + 1/2 \ \sigma + 1/2 \ 1/2) + (\sigma + 1/2 \ \sigma - 1/2 \ 1/2) + \cdots \\ &+ (\sigma + 1/2 \ 1/2 \ 1/2) + (\sigma - 1/2 \ \sigma - 1/2 \ 1/2) \\ &+ (\sigma - 1/2 \ \sigma - 3/2 \ 1/2) + \cdots + (\sigma - 1/2 \ 1/2 \ 1/2) + \cdots \\ &+ (3/2 \ 3/2) + (3/2 \ 1/2) + (1/2 \ 1/2), \end{aligned} \quad (2.3)$$

$$\begin{aligned} \sigma > \sigma' \quad (\sigma \ \sigma' \ 0) \otimes \left(\frac{1}{2} \frac{1}{2} \frac{1}{2} \right) &= \left(\sigma + \frac{1}{2} \ \sigma' + \frac{1}{2} \frac{1}{2} \right) \\ &+ \left(\sigma - \frac{1}{2}, \sigma' + \frac{1}{2} \frac{1}{2} \right) + \left(\sigma - \frac{1}{2} \ \sigma' - \frac{1}{2} \frac{1}{2} \right) \\ &+ (\sigma + 1/2 \ \sigma' - 1/2 \ 1/2), \end{aligned} \quad (2.4)$$

$$\begin{aligned} \sigma = \sigma' = \sigma \quad (\sigma \ \sigma \ 0) \otimes \left(\frac{1}{2} \frac{1}{2} \frac{1}{2} \right) &= (\sigma + 1/2 \ \sigma \\ &+ 1/2 \ 1/2) + (\sigma + 1/2 \ \sigma - 1/2 \ 1/2) \\ &+ (\sigma - 1/2 \ \sigma - 1/2 \ 1/2). \end{aligned} \quad (2.5)$$

这一约化是简单可约的, σ, σ' 都是正整数.

(3) Spin(6) 到 Spin(5) 的约化. 利用旋量与非旋量外积扬 (Young) 图法^[2], 我们导出下列约化规则

$$\begin{aligned} (\sigma_1 + 1/2 \ \sigma_2 + 1/2 \ 1/2) &= (\sigma_1 + 1/2 \ \sigma_2 + 1/2) + (\sigma_1 + 1/2 \ \sigma_2 - 1/2) \\ &+ (\sigma_1 + 1/2 \ \sigma_2 - 3/2) + \cdots + (\sigma_1 + 1/2 \ 1/2) + (\sigma_1 - 1/2 \ \sigma_2 \\ &+ 1/2) + (\sigma_1 - 1/2 \ \sigma_2 - 1/2) + (\sigma_1 - 1/2 \ \sigma_2 - 3/2) + \cdots \\ &+ (\sigma_1 - 1/2 \ 1/2) + (\sigma_1 - m + 1/2 \ \sigma_2 + 1/2) + (\sigma_1 - m \\ &+ 1/2 \ \sigma_2 - 1/2) + \cdots + (\sigma_1 - m + 1/2 \ 1/2). \end{aligned} \quad (2.6)$$

直到 $\sigma_1 - m = \sigma_2$ 时为止, σ_1, σ_2, m 都为正整数.

$$(\sigma + 1/2 \ \sigma + 1/2 \ 1/2) = (\sigma + 1/2 \ \sigma + 1/2) + (\sigma + 1/2 \ \sigma - 1/2)$$

表1 Spin(5)到

| τ | $(\tau \ 3/2)$ |
|--------|---|
| 3/2 | 9/2 5/2 3/2 |
| 5/2 | 13/2 11/2 9/2 (7/2) ² 5/2 3/2 1/2 |
| 7/2 | 17/2 15/2 (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² 3/2 1/2 |
| 9/2 | 21/2 19/2 (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² (3/2) ² |
| 11/2 | 25/2 23/2 (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² 3/2 1/2 |
| 13/2 | 29/2 27/2 (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² 3/2 1/2 |
| 15/2 | 33/2 31/2 (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² |
| 17/2 | 37/2 35/2 (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² |
| 19/2 | 41/2 39/2 (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² |
| 21/2 | 45/2 43/2 (41/2) ² (39/2) ² (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² |
| 23/2 | 49/2 47/2 (45/2) ² (43/2) ² (41/2) ² (39/2) ² (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² |
| 25/2 | 53/2 51/2 (49/2) ² (47/2) ² (45/2) ² (43/2) ² (41/2) ² (39/2) ² (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² |

| τ | $(\tau \ 5/2)$ |
|--------|---|
| 5/2 | 15/2 11/2 9/2 7/2 5/2 3/2 |
| 7/2 | 19/2 17/2 15/2 (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² 3/2 1/2 |
| 9/2 | 23/2 21/2 (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² (3/2) ² 1/2 |
| 11/2 | 27/2 25/2 (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² (3/2) ² 1/2 |
| 13/2 | 31/2 29/2 (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² (7/2) ² (5/2) ² |
| 15/2 | 35/2 33/2 (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² (11/2) ² (9/2) ² |
| 17/2 | 39/2 37/2 (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² (15/2) ² (13/2) ² |
| 19/2 | 43/2 41/2 (39/2) ² (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² (19/2) ² (17/2) ² |
| 21/2 | 47/2 45/2 (43/2) ² (41/2) ² (39/2) ² (37/2) ² (35/2) ² (33/2) ² (31/2) ² (29/2) ² (27/2) ² (25/2) ² (23/2) ² (21/2) ² |

$$+ \dots + (\sigma + 1/2 \ 1/2) \quad (2.7)$$

其中 σ 为正整数。这一约化是简单可约的。

(4) Spin(5) 到 Spin(3) 的约化。利用旋量与非旋量外积扬 (Young) 图法, 具体算出 Spin(5) 到 Spin(3) 的约化, 列于表 1。这一约化不是简单可约的。

三、波函数

我们解决了群链 (1.1) 的约化问题, 这一群链的 IBFM (II) 波函数可写成

$$|n_{\pi}, n_{\nu}\{n \ n'\} M(\Sigma_1 \Sigma_2 \Sigma_3)(\sigma_1 \sigma_2 \sigma_3)(\tau_1 \tau_2) K J M J\rangle, \quad (3.1)$$

其中 $\{n \ n'\}$ 标记 $U_{\pi}^{(6)}(6) \otimes U_{\nu}^{(6)}(6) = U_{\pi\nu}^{(6)}(6)$ 的不可约表示, 当 n_{π}, n_{ν} 给定时

$$\begin{aligned} \{n \ n'\} = & \{n_{\pi} + n_{\nu}\} + \{n_{\pi} + n_{\nu} - 1, 1\} + \{n_{\pi} + n_{\nu} - 2, 2\} \\ & + \dots + \{n_{\nu} \ n_{\pi}\}. \end{aligned} \quad (3.2)$$

其中 $n_{>} = \max(n_{\pi}, n_{\nu})$, $n_{<} = \min(n_{\pi}, n_{\nu})$.

$(\Sigma_1 \Sigma_2 \Sigma_3)$ 标记 $SO_{\pi}^{(6)}(6) \otimes SO_{\nu}^{(6)}(6) = SO_{\pi\nu}^{(6)}(6)$ 的不可约表示, $(\sigma_1 \sigma_2 \sigma_3)$ 标记 Spin

Spin(3) 的约化

| | 维数 |
|---|------|
| | 20 |
| | 64 |
| | 140 |
| | 256 |
| | 420 |
| | 640 |
| | 924 |
| $(5/2)^2 (3/2)^2$ | 1280 |
| $(9/2)^3 (7/2)^3 (5/2)^2 3/2 1/2$ | 1716 |
| $(13/2)^5 (11/2)^4 (9/2)^3 (7/2)^3 3/2 1/2$ | 2240 |
| $(17/2)^6 (15/2)^6 (13/2)^4 (11/2)^4 (9/2)^4 (7/2)^2 (5/2)^2 (3/2)^2$ | 2860 |
| $(21/2)^7 (19/2)^7 (17/2)^6 (15/2)^5 (13/2)^5 (11/2)^4 (9/2)^3 (7/2)^3 (5/2)^2 3/2 1/2$ | 3584 |
| $(25/2)^8 (23/2)^8 (21/2)^7 (19/2)^7 (17/2)^6 (15/2)^5 (13/2)^5 (11/2)^4 (9/2)^3 (7/2)^3 (5/2)^2 3/2 1/2$ | |

| | 维数 |
|--|------|
| | 56 |
| | 160 |
| | 324 |
| | 560 |
| | 880 |
| $(3/2)^2 1/2$ | 1296 |
| $(7/2)^4 (5/2)^3 (3/2)^2 1/2$ | 1820 |
| $(11/2)^6 (9/2)^5 (7/2)^4 (5/2)^3 (3/2)^2 1/2$ | 2464 |
| $(15/2)^7 (13/2)^7 (11/2)^6 (9/2)^5 (7/2)^4 (5/2)^3 (3/2)^2 1/2$ | 3240 |
| $(19/2)^8 (17/2)^8 (15/2)^7 (13/2)^7 (11/2)^6 (9/2)^5 (7/2)^4 (5/2)^3 (3/2)^2 1/2$ | |

(6)的不可约表示, $(\tau_1 \tau_2)$ 标记 Spin(5) 的不可约表示, K 是由于 Spin(5) 到 Spin(3) 的约化不是简单可约而引入的一个附加量子数。从上面的结果可以确定 (3.1) 式中各种量子数所能取的值。

四、动力学对称性

由于标记 IBFM(II) 波函数的群链是 (1.1) 式, 于是若原子核具有相应的动力学对称性, 即由 (1.1) 式群链标记的波函数是原子核的近似波函数, 则其哈密顿量可以写为

$$\begin{aligned}
 H = & a_{1\nu} C_{1U_{\nu}^{(B)}(6)} + a_{1\pi} C_{1U_{\pi}^{(B)}(6)} + b_1 C_{1U^{(F)}(4)} + a_2 C_{2U_{\pi+\nu}^{(B)}(6)} \\
 & + b_2 C_{2U^{(F)}(4)} + \epsilon_{2\pi} C_{1U^{(F)}(4)} C_{1U_{(\pi)}^{(B)}(6)} + \epsilon_{2\nu} C_{1U^{(F)}(4)} C_{1U_{\nu}^{(B)}(6)} \\
 & + \alpha C_{2SO_{\pi+\nu}^{(B)}(6)} + \beta C_{2Spin(6)} + \gamma C_{2Spin(5)} + \delta C_{2Spin(3)} \quad (4.1)
 \end{aligned}$$

其中 $C_{1U_{\nu}^{(B)}(6)}$, $C_{1U_{\pi}^{(B)}(6)}$, $C_{1U^{(F)}(4)}$ 和 $C_{1U_{\pi+\nu}^{(B)}(6)}$ 分别为 $U_{\pi}^{(B)}(6)$, $U_{\nu}^{(B)}(6)$, $U^{(F)}(4)$ 和 $U_{\nu+\pi}^{(B)}(6)$ 的一次 Casimir 算符, $C_{2U_{\pi+\nu}^{(B)}(6)}$, $C_{2U^{(F)}(4)}$, $C_{2SO_{\pi+\nu}^{(B)}(6)}$, $C_{2Spin(6)}$, $C_{2Spin(5)}$ 和 $C_{2Spin(3)}$ 分别为 $U_{\pi+\nu}^{(B)}(6)$, $U^{(F)}(4)$, $SO_{\pi+\nu}^{(B)}(6)$, Spin(6), Spin(5) 和 Spin(3) 的二次 Casimir 算符。和

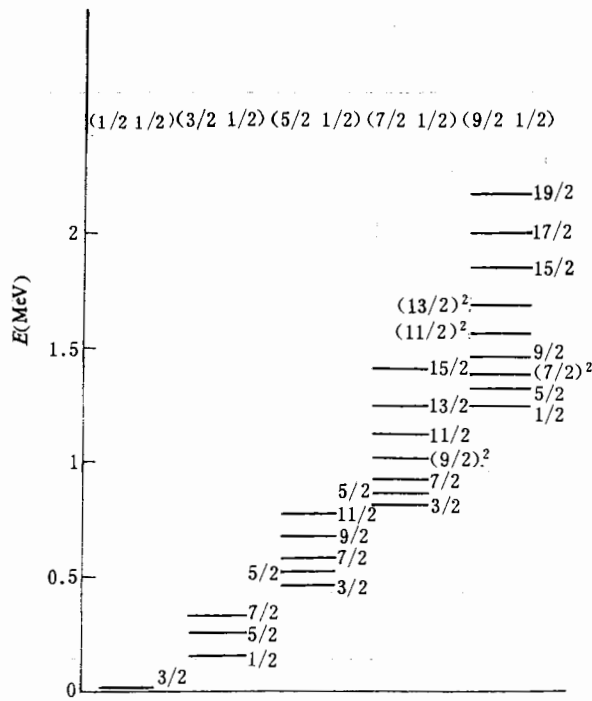


图1 是表示(17/2 1/2 1/2)所包含低于 2MeV 的能级

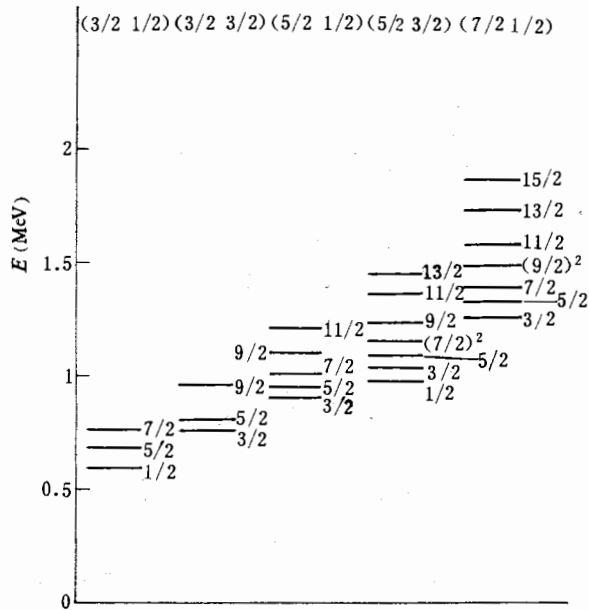


图2 是表示(15/2 3/2 1/2)所包含低于 2MeV 的能级

IBFM(II) 波函数 (3.1) 式对应的 H 的本征值为

$$E = a_{1\nu}n_\nu + a_{1\pi}n_\pi + b'_1M + b'_2M^2 + \varepsilon_{2\pi}Mn_\pi + \varepsilon_{2\nu}Mn_\nu + a_2[n(n+6) + n'(n'+18/5) - 2/5nn'] - \alpha/4[\Sigma_1(\Sigma_1+4) + \Sigma_2(\Sigma_2+2) + \Sigma_3^2]$$

$$- \beta/4[\sigma_1(\sigma_1 + 4) + \sigma_2(\sigma_2 + 2) + \sigma_3^2] + \gamma/6[\tau_1(\tau_1 + 3) + \tau_2(\tau_2 + 1)] + \delta J(J + 1) \quad (4.2)$$

其中 n_ν, n_π 分别是 ν 玻色子, π 玻色子的数目, M 是费米子数目, $a_{1\nu}, a_{1\pi}, b'_1, b'_2, \varepsilon_{2\pi}, \varepsilon_{2\nu}, a_2, \alpha, \beta, \gamma, \delta$ 都是可调参数.

IBFMII 在 Spin(6) 极限情况下典型能谱如图 1、2 所示.

五、结 论

(1) 从能谱中可以看出, 除 $U_{\pi\nu}^{(6)}(6)$ 的不可约表示 $\{n_\pi + n_\nu\}$ 包含的能级外, 还有 $\{n_\pi + n_\nu - 1, 1\}, \{n_\pi + n_\nu - 2, 2\} \cdots \{n_\pi + n_\nu - j, j\}$, 包含的能级, 换句话说, 能级丰富得多.

(2) 从典型能谱中可以看出, 即使对于 $\{n_\pi + n_\nu\}$ 的 IBFM (II) 的能级与 IBFM(I) IBM(I) 和 IBM(II) 的能级都有所不同.

3) 在 $M = 0$ 的全对称情况下和 $M = 1, j = 3/2$ 的情况下, 我们的所有结果均与文献[1]的结果一致. 我们的方法并不限于 $j = 3/2$. 当然 $j \neq 3/2$ 时其群链应作相应的变化.

4) 用上述方法我们业已对其它群链做了相应工作, 即将完成.

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THE SPIN (6) LIMIT OF INTERACTING BOSON-FERMION MODEL II

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ABSTRACT

The spin (6) limit of Interacting Boson-Fermion Model II is further discussed on the basis of the representation theory of Lie group. We generalize the Young-tableau methods for Kronecker products of group representations and discuss in detail the spinor symmetry of the IBFM II characterized by the group chain

$$U_{\pi}^{(B)}(6) \otimes U_{\nu}^{(B)}(6) \otimes U^{(F)}(4) \supset SO_{\pi}^{(B)}(6) \otimes SO_{\nu}^{(B)}(6) \otimes SU^{(F)}(4) \\ \supset SO_{\pi+\nu}^{(B)}(6) \otimes SU^{(F)}(4) \supset \text{Spin}(6) \supset \text{Spin}(5) \supset \text{Spin}(3)$$

We derive the reduction formulas in this group chain from our generalized Young-tableau methods. The wave functions of IBFM II are classified by this group chain. Some energy spectra are also calculated.